

Monte Carlo simulations of vortex dynamics in high-temperature superconductors under pulsed-field magnetization

Anna Moroz*, Anastasiia Maksimova, Vladimir Kashurnikov, Igor Rudnev

National Research Nuclear University MEPhI (Moscow Engineering Physics Institute), Moscow, Russia

*anmoroz@mephi.ru

ABSTRACT

Pulsed-field magnetization of a high-temperature Bi-2221 superconductor with various pinning landscapes has been studied by means of the Monte Carlo method. Computations have been performed for broad ranges of pulse parameters, such as amplitude and shape, and sample parameters, such as temperature and the concentration and type of distribution of defects (pinning centers). The time-dependences and distribution profiles of the average trapped magnetic field have been computed for various conditions. The differences between field-trapping efficiencies of different pinning landscapes have been demonstrated for various temperatures.

COMPUTATIONAL MODEL

$$G = \delta \left(\sum_{i < j} U_{\text{in-plane}}(r_{ij}) + \sum_i U_{\text{magn}}(x_i) + \sum_{i,j} U_{\text{bdr}}(|\mathbf{r}_i - \mathbf{r}_j^{\text{img}}|) + N\varepsilon + \sum_{i,j} U_{\text{pin}}(|\mathbf{r}_i - \mathbf{r}_j^{\text{def}}|) \right)$$

$$U_{\text{in-plane}}(r_{ij}) = 2U_0 K_0(r_{ij}/\lambda) \quad U_0 = (\Phi_0 / (4\pi\lambda))^2$$

$$\varepsilon = U_0 (\log(\lambda/\xi) + 0.52)$$

$$U_{\text{magn}}(x_i) = \lambda \sqrt{U_0} H \left(\frac{\cosh(x_i/\lambda)}{\cosh(L_x/2\lambda)} - 1 \right)$$

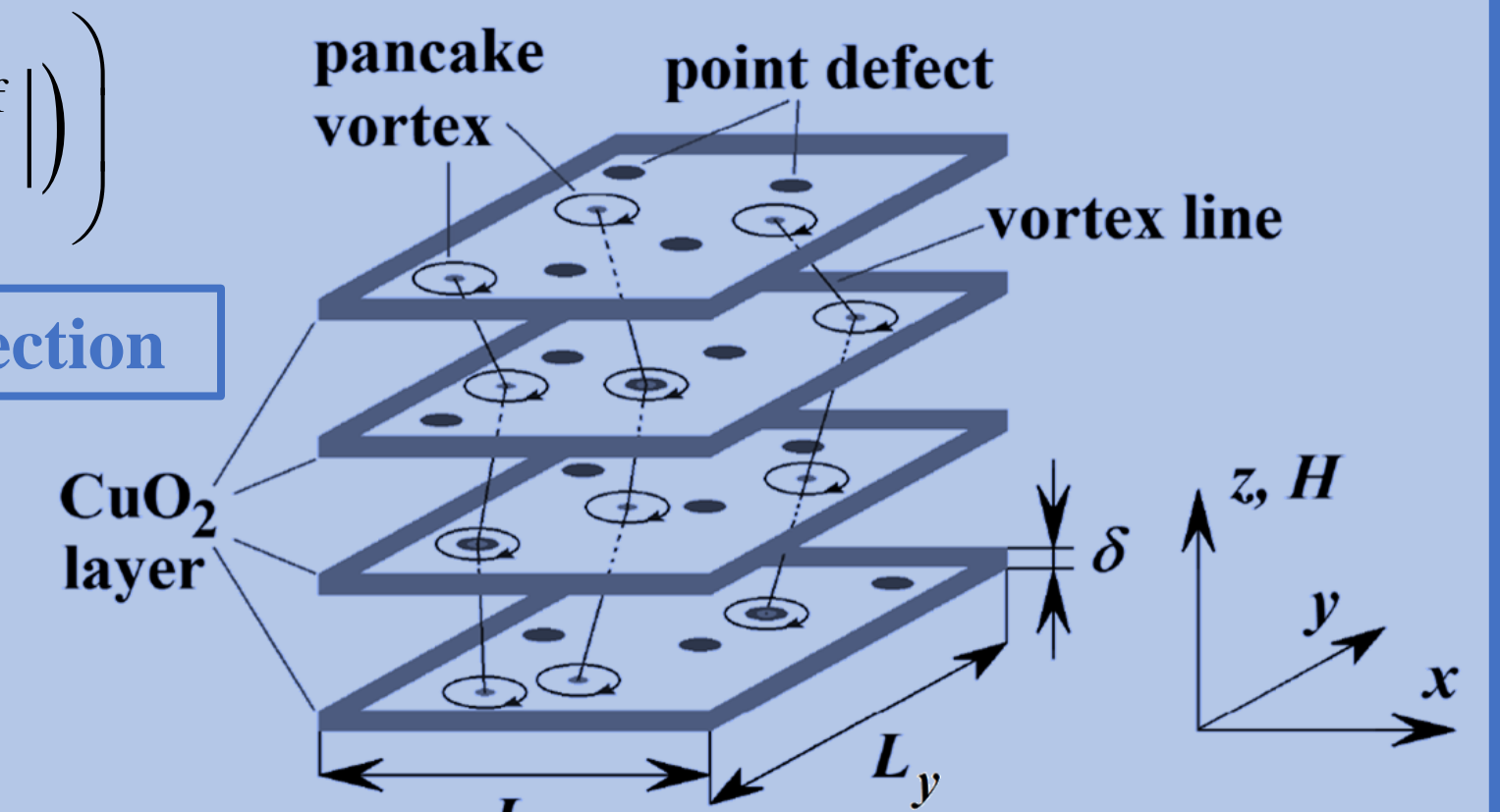
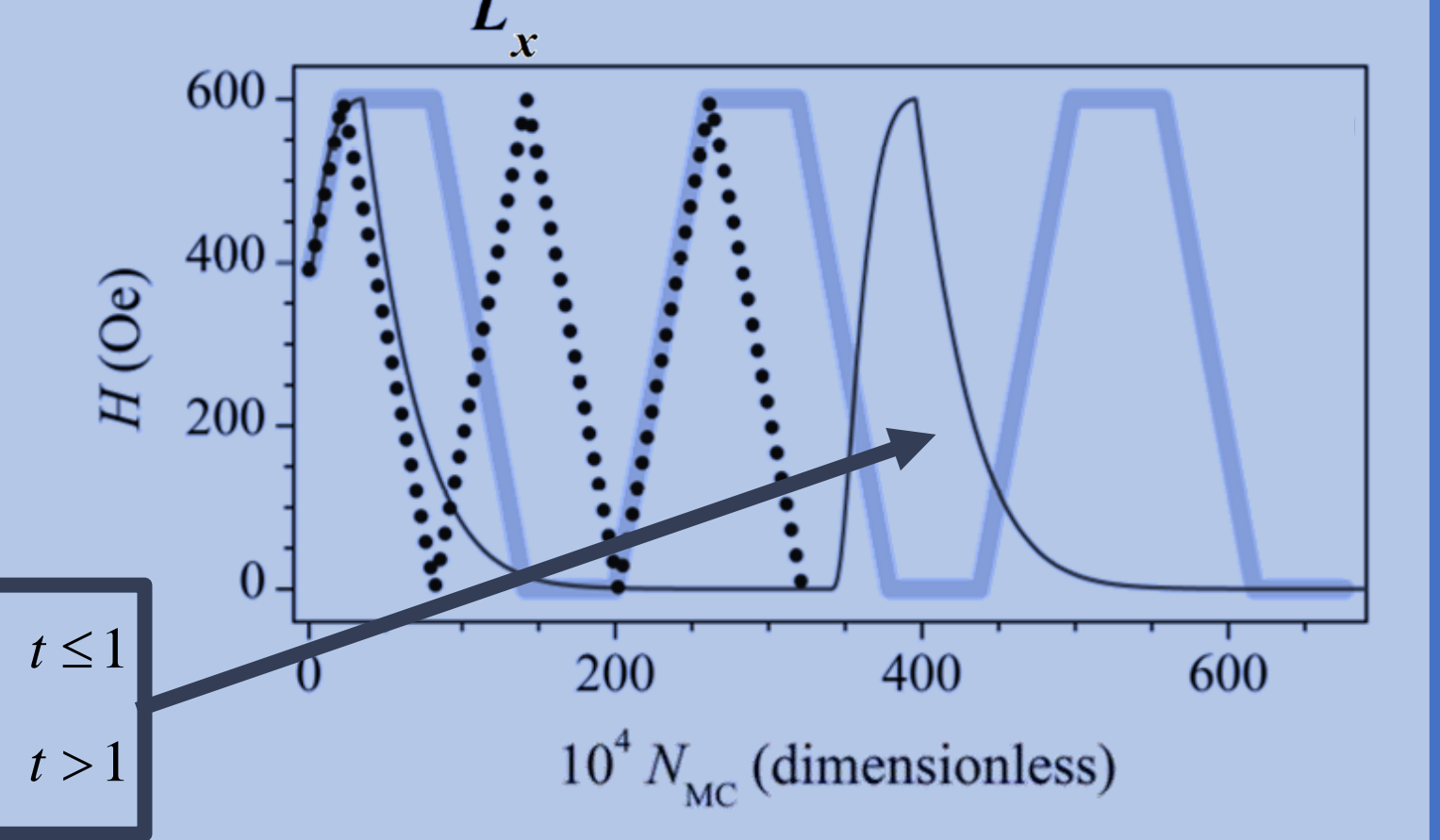
$$U_{\text{bdr}}(|\mathbf{r}_i - \mathbf{r}_j^{\text{img}}|) \equiv U_{\text{bdr}}(r_{ij}^{\text{img}}) = -U_0 K_0(r_{ij}^{\text{img}}/\lambda)$$

$$U_{\text{pin}}(|\mathbf{r}_i - \mathbf{r}_j^{\text{def}}|) \equiv U_{\text{pin}}(r_{ij}^{\text{def}}) = -\alpha \frac{\exp(-r_{ij}^{\text{def}}/2\xi)}{1 + r_{ij}^{\text{def}}/\xi}$$

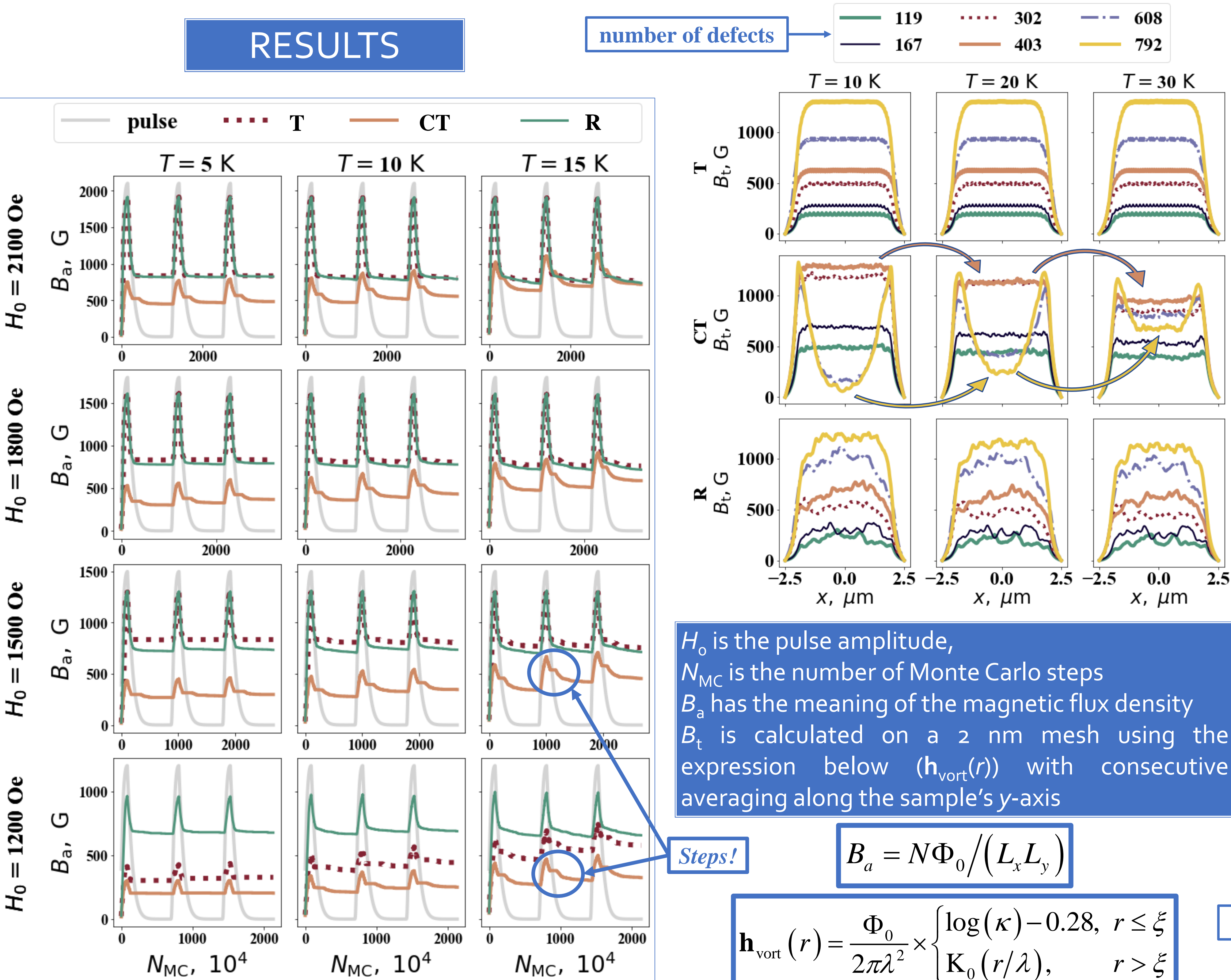
$$\Phi_0 = 2.07 \cdot 10^{-7} \text{ G cm}^2$$

$$H(t) = \begin{cases} H_0 t^{-1} \exp(1-t^{-1}), & t \leq 1 \\ H_0 \exp(1-t^{1.5}), & t > 1 \end{cases}$$

Bi₂Sr₂CaCu₂O_{8+x}
 $\lambda = 180 \text{ nm}, \xi = 2 \text{ nm}, T_c = 84 \text{ K}$
 $L_x = 5 \mu\text{m}, L_y = 3 \mu\text{m}, \delta = 0.27 \text{ nm}$

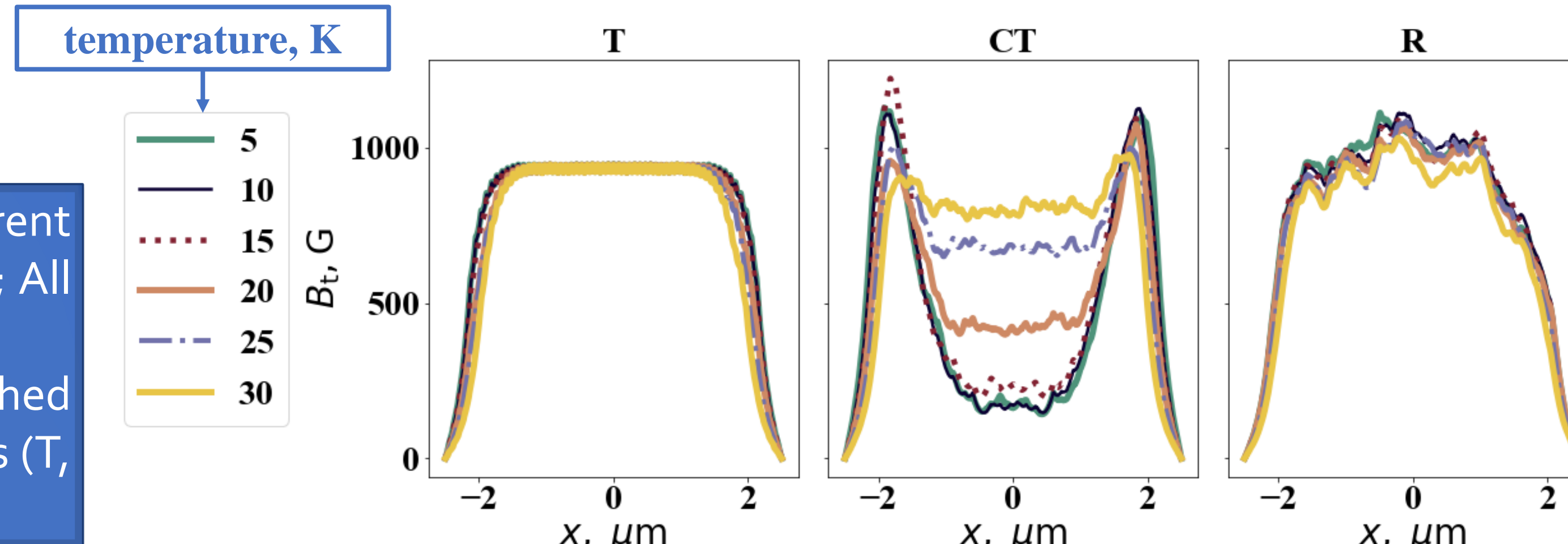



RESULTS



(on the left) The time-dependences of the averaged magnetic field induction B_a inside three samples with different pinning landscapes: a **triangular lattice (T)**, a **conformal triangular array (CT)**, and a **random distribution (R)**; All samples contain approximately the same number of defects (608) corresponding to a density of $4 \cdot 10^9 \text{ cm}^{-2}$;
 (on the right) The distribution profiles of the trapped magnetic field B_t (after the external field has been switched off) for samples with different numbers of defects (provided in the legend above) distributed in 3 different ways (T, CT and R as denoted on the left of the images); Here, $H_0 = 2100 \text{ Oe}$.

(on the right) The time-dependences of B_a inside samples with different pinning landscapes (T (top row), CT (central row), and R (bottom row) for three different pulse-shapes (trapezoidal (z), triangular (t) and exponential (e)). The image on the left and right correspond to different pulse amplitudes H_0 ; The samples contain 608 defects;
 (below) The distribution profiles of the trapped field inside samples at various temperatures (from 5 to 30 K). $H_0 = 2100 \text{ Oe}$



CONCLUSION

1) The CT array requires much higher pulse amplitudes for full magnetization but allows for more trapped flux at lesser defect densities;
 2) The increasing temperature causes thermoactivated flux creep which leads to successive magnetization upon repetitive application of pulses but also to step-like flux relaxation from CT arrays.